

mathematical methods - week 13

Probability

Georgia Tech PHYS-6124

Homework HW #13

due Monday, November 18, 2019

== show all your work for maximum credit,
== put labels, title, legends on any graphs
== acknowledge study group member, if collective effort
== if you are LaTeXing, here is the [source code](#)

Bonus points

Exercise **13.1** *Lyapunov equation*

12 points

This week there are no required exercises. Whatever you do, you get bonus points.

edited November 15, 2019

Week 13 syllabus

November 11, 2019

Mon A summary of key concepts

- [20.2 Moments, cumulants](#)

Wed Why Gaussians again?

- [33.2 Brownian diffusion](#)
- [33.3 Noisy trajectories](#)

Fri A glimpse of Ornstein-Uhlenbeck, the “harmonic oscillator” of the theory of stochastic processes. And the one “Lyapunov” thing Lyapunov actually did:)

- [Noise is your friend](#)
- [33.4 Noisy maps](#)
- 33.5 All nonlinear noise is local

13.1 Literature

Really going into the Ornstein-Uhlenbeck equation might take too much of your time, so this week we skip doing exercises, and if you are curious, and want to try your hand at solving exercise [13.1 Lyapunov equation](#), you probably should first skim through our lectures on the Ornstein-Uhlenbeck spectrum, Sect. 4.1 and Appen. B.1 [here](#). Finally! we get something one expects from a math methods course, an example of why orthogonal polynomials are useful, in this case the Hermite polynomials :)

The reason why I like this example is that again the standard ‘physics’ intuition misleads us. Brownian noise spreads with time as \sqrt{t} , but the diffusive dynamics of nonlinear flows is fundamentally different - instead of spreading, in the Ornstein-Uhlenbeck example the noise contained and balanced by the nonlinear dynamics.

- D. Lippolis and P. Cvitanović, *How well can one resolve the state space of a chaotic map?*, Phys. Rev. Lett. 104, 014101 (2010); [arXiv:0902.4269](#)
- P. Cvitanović and D. Lippolis, *Knowing when to stop: How noise frees us from determinism*, in M. Robnik and V.G. Romanovski, eds., *Let’s Face Chaos through Nonlinear Dynamics* (Am. Inst. of Phys., 2012); [arXiv:1206.5506](#)
- J. M. Heninger, D. Lippolis and P. Cvitanović, *Neighborhoods of periodic orbits and the stationary distribution of a noisy chaotic system*; [arXiv:1507.00462](#)

Question 13.1. Henriette Roux asks

Q What percentage score on problem sets is a passing grade?

A That might still change, but currently it looks like 60% is good enough to pass the course. 70% for C, 80% for B, 90% for A. Very roughly - will alert you if this changes. Here is the percentage score as of [week 10](#).

Question 13.2. Henriette Roux asks

Q How do I subscribe to the nonlinear and math physics and other seminars mailing lists?

A [click here](#)

Exercises

13.1. **Lyapunov equation.** Consider the following system of ordinary differential equations,

$$\dot{Q} = A Q + Q A^\top + \Delta, \quad (13.1)$$

in which $\{Q, A, \Delta\} = \{Q(t), A(t), \Delta(t)\}$ are $[d \times d]$ matrix functions of time t through their dependence on a deterministic trajectory, $A(t) = A(x(t))$, etc., with stability matrix A and noise covariance matrix Δ given, and density covariance matrix Q sought. The superscript $()^\top$ indicates the transpose of the matrix. Find the solution $Q(t)$, by taking the following steps:

- (a) Write the solution in the form $Q(t) = J(t)[Q(0) + W(t)]J^\top(t)$, with Jacobian matrix $J(t)$ satisfying

$$\dot{J}(t) = A(t) J(t), \quad J(0) = \mathbf{1}, \quad (13.2)$$

with $\mathbf{1}$ the $[d \times d]$ identity matrix. The Jacobian matrix at time t

$$J(t) = \hat{T} e^{\int_0^t d\tau A(\tau)}, \quad (13.3)$$

where \hat{T} denotes the ‘time-ordering’ operation, can be evaluated by integrating (13.2).

- (b) Show that $W(t)$ satisfies

$$\dot{W} = \frac{1}{J} \Delta \frac{1}{J^\top}, \quad W(0) = 0. \quad (13.4)$$

- (c) Integrate (13.1) to obtain

$$Q(t) = J(t) \left[Q(0) + \int_0^t d\tau \frac{1}{J(\tau)} \Delta(\tau) \frac{1}{J^\top(\tau)} \right] J^\top(t). \quad (13.5)$$

- (d) Show that if $A(t)$ commutes with itself throughout the interval $0 \leq \tau \leq t$ then the time-ordering operation is redundant, and we have the explicit solution $J(t) =$

$\exp \left\{ \int_0^t d\tau A(\tau) \right\}$. Show that in this case the solution reduces to

$$Q(t) = J(t) Q(0) J(t)^\top + \int_0^t d\tau' e^{\tau' A(t)} \Delta(\tau') e^{\tau' A^\top(t)}. \quad (13.6)$$

- (e) It is hard to imagine a time dependent $A(t) = A(x(t))$ that would be commuting. However, in the neighborhood of an equilibrium point x^* one can approximate the stability matrix with its time-independent linearization, $A = A(x^*)$. Show that in that case (13.3) reduces to

$$J(t) = e^{tA},$$

and (13.6) to what?